

Optical spectroscopy of X-Mega targets – V. The spectroscopic binary HD 93161 A and its visual companion HD 93161 B[★]

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ABSTRACT

We present the analysis of an extensive set of high-resolution spectroscopic observations of HD 93161, a visual binary with a separation of 2 arcsec. HD 93161 A is a spectroscopic binary, with both components clearly detected throughout the orbit. The primary star is most probably of spectral type O8V, while the secondary is likely an O9V. We obtain the first orbital solution for this system, characterized by a period of 8.566 ± 0.004 d. The minimum masses of the primary and secondary stars are 22.2 ± 0.6 and $17.0 \pm 0.4 M_{\odot}$, respectively. These values are quite large, suggesting a high inclination of the orbit. The second object, HD 93161 B, displays an O6.5V(f) spectral type and is thus slightly hotter than its neighbour. This star is at first sight single but presents radial velocity variations. Finally, we study HD 93161 in the X-ray domain. No significant variability is detected. The X-ray spectrum is well described by a 2T model with $kT_1 \sim 0.3$ keV and $kT_2 \sim 0.7$ keV. The X-ray luminosity is rather moderate, without any large emission excess imputable to a wind interaction.

Key words: binaries: spectroscopic – stars: early-type – stars: fundamental parameters – stars: individual: HD 93161 (A,B) – X-rays: individual: HD 93161 – X-rays: stars.

1 INTRODUCTION

A few years ago began an extensive campaign to study the Carina nebula, one of the largest concentrations of young and massive stars known in our Galaxy. This campaign was initiated by the X-Mega group (Corcoran et al. 1999) in order to gain a good knowledge of the cluster's binary population prior to its observation in the X-ray domain by the *ROSAT* satellite. Such multiwavelength observations of spectroscopic binaries in open stellar clusters enable us to constrain the fundamental parameters of early-type stars and to investigate the interactions between their components in more depth.

The efforts were focused more specifically on the massive stars of the Trumpler 16 (Tr 16) open cluster (Albacete Colombo et al. 2001; Morrell et al. 2001; Rauw et al. 2001; Albacete Colombo et al.

2002, hereafter Papers I–IV, respectively). This cluster constitutes one of the best places to complete our objective, because it harbours a large number of binaries; Levato et al. (1991) found at least five spectroscopic binaries and several radial velocity (RV) variables in Tr 16. The (very) preliminary results of Levato et al. (1991), relying on poor-quality and scarce data, were subsequently revised in the first papers of this series.

In the present paper, we analyse an extensive set of spectra of one of the still poorly studied systems of Tr 16, HD 93161 (= Tr 16–176, CPD–58°2631).¹ The compound formed by HD 93161 and HD 93160 actually harbours five stellar components, noted A–E, some of which were found by speckle observations (Mason et al. 1998): HD 93160 corresponds to the C component, located at 13 arcsec from the AB pair; component D is at 3 arcsec from HD 93160 and component E is at 8 arcsec of the AB pair. The latter two stars, separated by only 2 arcsec, have been known for a long time (see, for example, Walborn 1973) and are generally referred to as HD 93161. Penny (1996) and Howarth et al. (1997) reported RV variations in the ultraviolet (UV) spectrum of HD 93161. Unfortunately, the small separation between A and B, together with the low resolution of the *International Ultraviolet Explorer* (*IUE*), have not permitted us to attribute these changes confidently to one member of the pair. Howarth et al. (1997) even suggested that

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¹ The membership of this system to Tr 14, Tr 16 or Collinder 232 has been subject to controversy throughout the literature. Here we follow Walborn (1973) in considering that it belongs to Tr 16.

both stars could be short-period binaries. Previously, Levato et al. (1991) had provided an SB1 orbital solution for HD 93161, without specifying to which component this solution applied. A thorough study of the system was thus eagerly awaited.

This paper is organized as follows. The observations and the data reduction are presented in Section 2 and the spectral features of HD 93161 A and B are discussed in Sections 3 and 4. In Section 5, we derive the first full orbital solution for HD 93161 A, whereas Section 6 deals with the analysis of its X-ray emission. As a next step, we examine the evolutionary status of the stars and we finally conclude in Section 8.

2 OBSERVATIONS AND DATA REDUCTION

2.1 Optical spectroscopy

HD 93161 lies in a crowded region of the Carina nebula, with HD 93160 just 13 arcsec west of the star. As noted in the introduction, HD 93161 itself was found to be a visual double system, composed of two stars with similar brightnesses separated by only 2 arcsec. As in Vreux & Andrillat (1974), we define component A as being the one closest to HD 93160. With such a small separation between A and B, it is difficult to study these stars individually. During some observing nights, a poor seeing and/or the low angular resolution of the pointing camera did not enable us to clearly disentangle the close pair AB. Among our 45 spectra (see Table 1), only three can securely be attributed to the sole B component. The main contributor to the remaining 42 spectra is intended to be the A component, but these observations are usually contaminated, to a variable degree, by the B companion. This is marked in the spectra by the presence of a third, weak component.

We have observed HD 93161 over a period of 7 yr with various instruments. A first set of six medium-resolution spectra covering the wavelength range 3850–4800 Å was gathered in 1997 with the European Southern Observatory (ESO) 1.5-m telescope equipped with a Boller & Chivens (hereafter B&C) Cassegrain spectrograph. The data were obtained with a holographic grating (2400 lines mm⁻¹, ESO grating 32) providing a reciprocal dispersion of 32.6 Å mm⁻¹. The detector was a thinned, UV flooded Loral-Lesser CCD (ESO 39). The slit width was set to 220 µm corresponding to 2 arcsec on the sky. The spectral resolution, as measured on the helium–argon calibration spectra, is 1.2 Å, corresponding to a resolving power of 3600. Typical exposure times were of the order of 10 min and the average signal-to-noise (S/N) ratio was about 250. Parts of the B and C spectra were affected by a fringing pattern (Turatto, Tighe & Castillo 1997) that occurred over the wavelength range 4050–4250 Å in this instrument configuration. Given the variability of the fringing pattern and in order to avoid amplification of the fringes in the stellar spectra, the data were not flat-fielded. All the reductions were performed using the MIDAS software developed at the ESO. Note that only B and C spectra were sky-subtracted.

Six high-resolution spectra of HD 93161 were taken in 1999 May–June with the bench-mounted echelle spectrograph (BME) fed by the 1.5-m Ritchey–Chrétien telescope at the CTIO. These data covered the spectral range 3750–5800 Å. The typical S/N ratio was ~60 at 5430 Å for exposure times of 1 h. The detector used was a Tek CCD with 2048 × 2048 pixels of 24 × 24 µm² and the resolving power was 45 000. The BME data were reduced using the IRAF² package, following the recommendations of the BME User's Manual. A first rectification of the extracted orders was carried out

with the projector flat exposures. The spectra were then normalized by fitting a low-order polynomial to the continuum.

Between 1999 and 2004, we obtained 33 additional high-resolution spectra at the La Silla Observatory with the Fibre-fed Extended Range Optical Spectrograph (FEROS; Kaufer et al. 1999). This spectrograph was installed until 2002 October at the ESO 1.5-m telescope, and was subsequently moved to the 2.2-m telescope. The FEROS spectra cover the wavelength range 3750–9000 Å with a resolving power of 48 000. The detector was a 2 k × 4 k EEV CCD with pixel size 15 × 15 µm². The typical exposure time was 10–15 min and the S/N ratio at 5650 Å is about 125–225. The FEROS spectra were reduced using the appropriate context of the MIDAS environment, together with an improved reduction pipeline (Sana et al. 2003). They were normalized by fitting the continuum with a polynomial of degree ranging from 2 to 6. Most of the unavoidable fringes affecting the red part of the spectra were corrected by flat-fielding, but there are a few residuals, e.g. near 6640 Å.

2.2 X-ray data

The core of the Carina nebula was observed several times by the *XMM-Newton* observatory (Jansen et al. 2001). We focus here on the five data sets acquired in 2000 and 2001. The first two data sets, centred on η Carinae, were obtained in 2000 July during satellite revolutions 115 and 116. The three other observations, centred on WR 25, were taken 1 yr later during revolutions 283, 284 and 285. A preliminary study using only the first two observations was presented by Albacete Colombo, Méndez & Morrell (2003). A thorough analysis of the field relying on the entire observational data set has been performed by Antokhin et al. (in preparation).

HD 93161 was observed serendipitously during all five *XMM-Newton* pointings (see Table 2, later). The source was recorded with the three European Photon Imaging Camera (EPIC) detectors (PN, MOS1 and MOS2; for more details, see Strüder et al. 2001; Turner et al. 2001). The data were processed with the *XMM-Newton* Science Analysis Software (SAS, version 5.4.1). A check on proton flares resulted in the deletion of part of the exposure time. More details on the pipeline processing of these data can be found in Raassen et al. (2003).³

3 SPECTRUM OF HD 93161 A

Up to now, the visible spectrum of HD 93161 was essentially studied through photographic observations. Following Walborn (1973), both stars of HD 93161 have identical spectra and apparent magnitudes. Walborn further attributes an O6.5V(f) type to each star. The following year, however, Vreux & Andrillat (1974) discovered that HD 93161 A was cooler than HD 93161 B. Unfortunately, no formal spectral type was assigned to HD 93161 A by Vreux & Andrillat (1974).

We have investigated the spectrum of HD 93161 A through high-resolution echelle spectroscopy. Such a spectrum is shown in Fig. 1; the lines are clearly double, showing that this star is indeed a spectroscopic binary. Beside the lines of the Balmer series, the spectrum reveals many absorption lines of He I, He II, N III, Si III, Si IV, C IV and O III that are typical of O stars. It also reveals a number of diffuse interstellar bands (DIBs), as well as interstellar absorptions due to He I λ3889, Na I, Ca I, Ca II, CH and CH⁺. There are only few emission lines in the spectrum of HD 93161 A. Most of them

² IRAF is distributed by the NOAO.

³ In Raassen et al. (2003), note that column 4 of Table 5 actually gives the Julian Dates in the format JD–245 0000, not the MJDs.

Table 1. Journal of the spectroscopic observations of HD 93161 A and B. Column 1 gives the instrument used and column 2 lists the Heliocentric Julian Date (HJD) at mid-exposure. The next three columns present the phases (calculated from the He I $\lambda 4471$ orbital solution with $e = 0$) and the He I $\lambda 4471$ RVs of the primary and secondary components. The last three columns are similar, but refer to the averaged velocities (see Section 5).

Instrument	HJD–245 0000 (d)	ϕ_{4471}	RV ₁ (km s ^{−1})	RV ₂	ϕ_{avg}	$\overline{\text{RV}}_1$ (km s ^{−1})	$\overline{\text{RV}}_2$
HD 93161 A							
B&C	534.752	0.085	90.0	−105.4			
	535.758 ^a						
	536.748	0.318	111.8	−167.0			
	537.741	0.434	81.8	−83.9			
	538.733	0.550	−20.1	−20.1			
BME	539.661	0.658	−151.0	100.9			
	1328.625	0.760	−178.5	194.6	0.761	−156.9	173.5
	1329.556	0.869	−112.2	127.7	0.870	−117.3	125.8
	1330.590	0.990	−6.7	−6.7	0.991	−5.5	−4.4
	1331.674	0.116	81.1	−141.9	0.117	91.3	−130.2
FEROS	1332.560	0.220	149.6	−195.4	0.221	153.3	−192.6
	1299.554	0.367	102.8	−144.2	0.368	111.1	−150.4
	1300.563	0.484	−4.0	−4.0	0.485	−5.9	−5.0
	1301.670	0.614	−100.9	131.6	0.615	−102.1	130.8
	1302.662	0.729	−157.6	204.6	0.731	−150.3	198.2
	1304.560	0.951	−54.7	66.6	0.952	−44.7	62.5
	1327.531	0.633	−98.9	155.1	0.634	−102.6	151.9
	1669.577	0.562	−55.6	81.9	0.563	−58.4	79.8
	1670.539	0.675	−139.7	184.2	0.675	−143.3	184.2
	1671.535	0.791	−151.8	195.7	0.791	−149.2	192.4
	1672.527	0.907	−92.7	121.0	0.907	−95.7	109.7
	1673.526	0.023	−3.6	−3.6	0.024	−2.0	−1.1
	2037.627	0.528	−70.7	19.8	0.528	−7.9	18.6
	2038.556	0.636	−126.1	107.5	0.636	−120.5	112.4
	2039.628	0.761	−159.8	211.4	0.761	−156.5	205.7
	2040.636	0.879	−79.0	133.4	0.879	−93.0	133.5
	2335.615	0.314	112.6	−192.8	0.314	118.4	−191.0
	2336.617	0.431	24.7	−102.4	0.431	52.7	−72.6
	2337.606	0.547	−70.3	80.9	0.546	−64.7	49.5
	2338.604	0.663	−111.7	173.4	0.663	−114.5	167.5
	2381.573	0.679	−138.5	185.4	0.679	−139.4	182.4
	2382.552	0.794	−156.8	195.8	0.793	−154.3	190.8
	2383.556	0.911	−75.7	105.2	0.910	−79.1	103.4
	2782.527	0.486	−2.2	−2.2	0.485	−2.4	−0.9
	2783.516	0.601	−84.2	121.5	0.600	−87.9	121.0
	2784.511	0.717	−157.4	205.3	0.716	−154.4	200.8
	3130.515	0.109	101.2	−118.2	0.108	100.7	−122.4
	3131.495	0.224	161.2	−192.8	0.222	160.8	−193.3
	3132.494	0.340	141.9	−170.4	0.339	141.9	−170.8
	3133.575	0.466	39.9	−55.3	0.465	39.7	−50.0
	3134.513	0.576	−65.4	97.7	0.574	−68.1	94.4
	3135.544	0.696	−146.9	195.2	0.695	−147.1	193.1
HD 93161 B							
BME	1330.640		−4.5				
FEROS	2382.563		−57.8				
	3132.506		−14.0				

^aNo measurement is available for He I $\lambda 4471$, but the He II $\lambda 4686$ line shows RVs of 109.2 and -161.2 km s^{−1} for the primary and secondary, respectively.

are sharp nebular lines, associated with the Carina nebula. We may note that the most massive component exhibits C III $\lambda 5696$ in emission.⁴ Throughout this paper, we refer to the more massive star as the primary, or HD 93161 A1.

⁴ The fact that only this star exhibits the C III $\lambda 5696$ line in emission helps for the identification of the primary star. In addition, the presence of a second emission peak in this line clearly indicates pollution by HD 93161 B, which

Only the least blended and higher-quality spectra are suitable for the spectral classification. Following Conti & Alschuler (1971); Conti & Frost (1977) and Mathys (1988, 1989), we adopt the log [EW(4471)/EW(4542)] classification criterion to determine the

also displays this line in its spectrum. It thus enables us to estimate the degree of contamination by HD 93161 B in each spectrum of HD 93161 A.

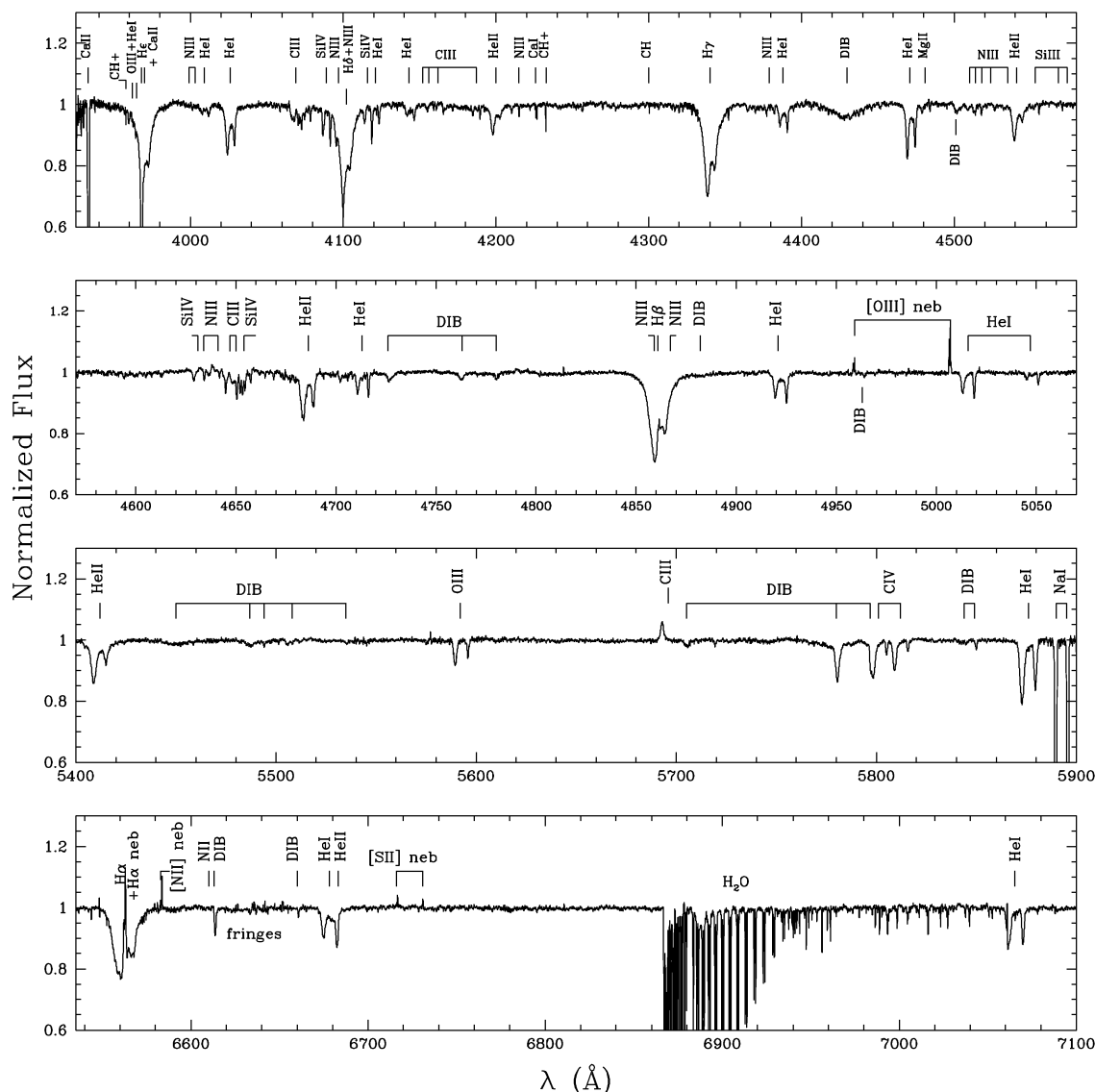


Figure 1. Spectrum of HD 93161 A as observed with the FEROS spectrograph on HJD 245 3135.544 ($\phi \sim 0.7$). At this orbital phase, the spectrum of the most massive star is shifted towards the blue. The main lines are indicated at their rest wavelengths.

spectral types of the stars. For this purpose, we fitted these lines with a combination of several Gaussians. This procedure enables us to measure the RVs rather securely, but the presence of small additional components (due to the contamination by the neighbouring HD 93161 B) renders the determination of accurate equivalent widths (EWs) rather awkward. The measured ratios are thus rather dispersed, resulting in an approximate O7–7.5 type for the primary and an O7–8.5 type for the secondary at 1σ dispersion. To refine this classification further, we have compared their spectrum to the OB star atlas of Walborn & Fitzpatrick (1990). For the primary star, He I $\lambda 4471$ (resp. 4026) appears only slightly stronger than He II $\lambda 4542$ (resp. 4200). The reverse situation occurs for the He I $\lambda \lambda 4143, 4388$ lines, which are weaker than the He II $\lambda \lambda 4200, 4542$ lines. The spectrum of HD 93161 A1 is thus similar to that of HD 48279 (Walborn & Fitzpatrick 1990), indicating an O8 type for this star. On the other hand, the secondary seems cooler; He I $\lambda 4471$ (resp. 4026) is much stronger than He II $\lambda 4542$ (resp. 4200), while He I $\lambda 4143$ is weaker than He I $\lambda 4388$ but has a strength similar to those of He II $\lambda \lambda 4200,$

4542. HD 93161 A2 thus resembles HD 46202, an O9 star. These spectral types are slightly later than those derived from Conti's criterion. This difference is probably due to the contamination by the B component but we cannot totally exclude that the stars might be one subtype earlier.

For the luminosity classification, we rely on the $\log[\text{EW}(\text{Si IV } \lambda 4089)/\text{EW}(\text{He I } \lambda 4143)]$ criterion introduced by Conti & Alschuler (1971). For the same reason as before, the measured values of this ratio are rather scattered but they clearly indicate a main-sequence classification. Note that the $\log[\text{EW}(\text{He II } \lambda 4686)] + \log[\text{EW}(\text{He I } \lambda 4388)]$ luminosity criterion introduced by Mathys (1988, 1989) was not used here; it relies on absolute EWs, which are difficult to determine because of the variable contamination by HD 93161 B. The main-sequence classification is further supported by the comparison of our spectra with the atlas of standards because, for both stars, Si IV $\lambda 4089$ is quite faint, whereas the He II $\lambda 4686$ line appears rather strongly in absorption. We thus conclude that HD 93161 A is most probably an O8V+O9V binary.

hamper the RV determination. These RV variations might actually be caused by orbital motion in a binary. Unfortunately, the small number of B spectra does not enable us to further analyse the variability of this star, and especially to determine if these variations are periodic. Additional uncontaminated, high S/N data of HD 93161 B would certainly allow us to investigate the multiplicity of this star.

5 ORBITAL SOLUTION FOR HD 93161 A

As mentioned before, we have determined the RVs of the lines by fitting them with a sum of Gaussian profiles. In order to obtain a reliable orbital solution, we considered 12 different lines: He I $\lambda\lambda 4026, 4471, 4921, 5016, 5876, 7065$; He II $\lambda\lambda 4542, 4686, 5412$; O III $\lambda 5592$ and C IV $\lambda\lambda 5801, 5812$. We present in Table 1 the RVs measured for the He I $\lambda 4471$ line. Note that we have also measured the RVs of the most prominent interstellar lines; for our FEROS observations, the average RV of the CH⁺ $\lambda 4232$ line is $5.3 \pm 0.4 \text{ km s}^{-1}$, whereas for the B and C data, the mean RV of the Ca II $\lambda 3933$ line is $-35 \pm 15 \text{ km s}^{-1}$ and for the BME data, the mean RV of the red component of the Ca II $\lambda 3933$ line is $-0.1 \pm 2.9 \text{ km s}^{-1}$. We thus emphasize the very good temporal stability of the FEROS spectrograph, whose spectra constitute the majority of our sample.

Up to now, two orbital solutions for HD 93161 have been published in the literature. Levato et al. (1991) have calculated an SB1 solution for HD 93161, but without specifying to which component (either A or B) it refers. Luna et al. (2003) failed to provide any improvement to the results of Levato et al. (1991).⁵ If these authors measured the RVs on combined spectra of HD 93161, as can be guessed from their description, it seems rather strange that they detected only one component (thus deriving an SB1 solution) because the presence of multiple components is easily spotted even in our medium-resolution B and C spectra. In the rest of this paper, we do not discuss these analyses further.

5.1 Period determination

Once the RVs are measured, the next step is to determine the period of the orbital motion. To this aim, we have applied the method of Lafler & Kinman (1965, hereafter LK) as well as a Fourier-type analysis (the method of Heck, Manfroid & Mersch 1985; see also the comments by Gosset et al. 2001). Because confusion could be a possible issue, these period search algorithms were first applied to the time series of the absolute RV differences $|RV_1 - RV_2|$; for each of the lines quoted above, a period of about 4.3 d was then detected. Using this preliminary period and the RVs of the C III $\lambda 5696$ line, we then identified the primary and secondary, and applied the period search on the actual differences $RV_1 - RV_2$. Both the Fourier analysis and the LK algorithm agree on a value of $P = 8.565 \text{ d}$ for all lines.

Our data set consists of 42 spectra spread over 2600 d for the He I $\lambda 4471$ and He II $\lambda 4686$ lines (where there exist B and C data in addition to the FEROS and BME observations) and of 36 spectra

spread over 1836 d for the other lines. These observational time bases correspond to natural peak widths $\Delta\nu$ of 0.00038 and 0.00054 d⁻¹, respectively. The full width at half-maximum of the actual peaks in the periodogram is in agreement with these theoretical values. We adopt an uncertainty of one-tenth of the peak's width (i.e. $P = 8.565 \pm 0.003$ or 0.004 d), the uncertainty depending on whether or not the B and C data are considered.

5.2 Orbital elements

To derive the orbital elements, we used a modified version of the Wolfe, Horak & Storer (1967) method. The algorithm, presented by Sana et al. (2003), allows a weighting that differs for the secondary and primary RVs. In the present paper, we have adopted a period value and a relative (primary versus secondary) weight that yield the lower residuals. In addition, we accounted for the quality of the measurements by attributing lower weights to the observations where the components are blended (i.e. only one measurement is available for both stars), noisy (e.g. BME and B and C data), or heavily contaminated by HD 93161 B.

At first, we fitted the RVs from each line with a general, eccentric solution. However, the computed eccentricity e turned out to be very small (<0.03), with a null eccentricity within less than 2σ . Therefore, a circular orbit was assumed for HD 93161 A. The results of these non-eccentric fits are displayed in Table 3 for each line. They are compatible with each other, within the errors.

Finally, we tried to obtain more accurate results by averaging the RVs from different lines. This was done after subtracting the apparent systemic velocities derived for each line from the individual fits. Note that the lines yielding the largest residuals (i.e. He II $\lambda\lambda 4542, 5412$ and He I $\lambda 4026$) were not considered. The lower-resolution B and C data were also discarded in order to have a more homogeneous set of RVs. Averaging all other observations allows us to obtain an rms residual of 8 km s^{-1} . This best solution, including a reappraisal of the period ($P_{\text{final}} = 8.5663 \pm 0.0040 \text{ d}$), is presented in Table 4 and is shown in Fig. 3; it is adopted throughout the rest of the paper.

The minimum masses $m \sin^3 i$ derived for the primary and secondary are 22.2 and 17.0 M_{\odot} , respectively (see Table 4). Comparing these values to typical masses of O8V and O9V stars allows us to estimate the inclination of the system. However, one must be careful when doing so, because the analysis of the binaries in the Tr 16 cluster (HD 93205, Paper II; Tr 16–104, Paper III) has shown that the typical masses presented by well-known references (such as, for example, Howarth & Prinja 1989) are generally too large

Table 2. Journal of the *XMM-Newton* observations of HD 93161. The first column gives the revolution during which the observations were taken, column 2 the Julian Date (JD) at mid-exposure, column 3 the phases (calculated from the average He I orbital solution; see Section 5), and the last three columns the exposure times in ks for each instrument. For more details, see Raassen et al. (2003).

Rev.	JD–245 0000 (d)	ϕ_{avg}	Exposure time (ks)		
			MOS1	MOS2	PN
115	1751.92	0.18	33.8	30.6	31.7
116	1753.56	0.37	11.2	8.3	9.4
283	2086.00	0.17	36.7	36.7	34.6
284	2089.05	0.53	42.1	42.1	39.6
285	2090.91	0.75	37.4	37.4	34.9

⁵ Note that, in Luna et al. (2003), the figure showing the ‘new solution’ for HD 93161 is most probably erroneous. Moreover, in the same bulletin, the authors also provide an SB1 solution for Tr 16–104 and claim to detect no RV variations for Tr 16–110, although both stars were already known as triple systems prior to this publication (Papers III and IV). This may legitimately cast some doubt on the quality of their data.

Table 3. Circular orbital solutions for HD 93161 A as deduced from different lines. K_1 and K_2 yield the amplitude of the RV curves of the primary and secondary, respectively, whilst γ_1 and γ_2 correspond to the apparent systemic velocities of the two components. To account for a possible velocity gradient in the line formation region, we allowed γ_1 and γ_2 to be different. Quoted uncertainties are the 1σ error bars. s_y/s_x gives the ratio of the secondary to primary uncertainties on the RVs.

Lines	s_y/s_x	P (d)	K_1 (km s $^{-1}$)	K_2 (km s $^{-1}$)	γ_1 (km s $^{-1}$)	γ_2 (km s $^{-1}$)	rms (km s $^{-1}$)
He I λ 4026	0.8	8.5665	149.7 \pm 3.4	201.2 \pm 4.6	5.1 \pm 4.0	13.0 \pm 5.0	14.8
He I λ 4471	1.0	8.5662	151.5 \pm 2.4	202.2 \pm 3.2	−1.4 \pm 2.5	2.2 \pm 3.1	12.4
He II λ 4542	1.4	8.5664	150.8 \pm 4.8	198.1 \pm 6.4	4.7 \pm 7.7	2.9 \pm 9.6	20.1
He II λ 4686	1.1	8.5663	150.8 \pm 2.4	201.1 \pm 3.2	9.1 \pm 2.5	3.0 \pm 3.0	12.4
He I λ 4921	0.8	8.5662	148.9 \pm 2.3	202.3 \pm 3.1	−1.1 \pm 2.1	1.5 \pm 2.4	9.4
He I λ 5016	0.8	8.5662	155.3 \pm 1.6	201.0 \pm 2.1	0.6 \pm 1.3	−0.2 \pm 1.4	6.6
He II λ 5412	1.3	8.5666	150.8 \pm 4.1	192.9 \pm 5.2	9.0 \pm 7.5	−0.2 \pm 9.5	20.8
O III λ 5592	1.2	8.5664	155.2 \pm 2.3	201.3 \pm 3.0	0.7 \pm 2.3	−4.2 \pm 2.7	10.4
C IV λ 5801	1.2	8.5663	153.0 \pm 2.7	198.3 \pm 3.5	−2.5 \pm 2.9	−5.4 \pm 3.5	11.3
C IV λ 5812	1.4	8.5663	155.1 \pm 2.4	197.9 \pm 3.1	−5.1 \pm 2.5	−2.8 \pm 3.0	10.0
He I λ 5876	0.9	8.5663	152.5 \pm 2.4	202.2 \pm 3.2	3.9 \pm 2.4	3.1 \pm 2.8	9.8
He I λ 7065	0.9	8.5663	155.8 \pm 2.2	201.4 \pm 2.9	2.9 \pm 2.0	4.1 \pm 2.2	8.4

Table 4. Orbital and physical parameters of HD 93161 A as deduced from the mean RVs for a circular orbit. The usual notations are used, R_{RL} being the Roche lobe radii of the stars and T_0 corresponding to the conjunction with the primary star in front. The error on the period is calculated from the observational time base (see Section 5.1) and all other errors are estimated 1σ deviations calculated for a fixed s_y/s_x ratio and a fixed period (those that give the lowest residuals). Note that the apparent systemic velocities (γ_1 and γ_2) of each line were subtracted prior to averaging.

P	8.5663 \pm 0.0040 d
T_0 (HJD −245 0000)	3001.098 \pm 0.011
s_y/s_x	0.9
m_1/m_2	1.31 \pm 0.02
K_1	152.9 \pm 1.7 km s $^{-1}$
K_2	200.6 \pm 2.2 km s $^{-1}$
γ_1 (after sub.)	0.02 \pm 1.4 km s $^{-1}$
γ_2 (after sub.)	0.05 \pm 1.6 km s $^{-1}$
$a_1 \sin i$	25.9 \pm 0.2 R_\odot
$a_2 \sin i$	33.9 \pm 0.4 R_\odot
$m_1 \sin^3 i$	22.2 \pm 0.6 M_\odot
$m_2 \sin^3 i$	17.0 \pm 0.4 M_\odot
$R_{RL}^1/(a_1 + a_2)$	0.4026 \pm 0.0015
$R_{RL}^2/(a_1 + a_2)$	0.3557 \pm 0.0014
rms	7.7 km s $^{-1}$

compared to the masses derived from the study of binary systems. We therefore decide to adopt the values determined for detached eclipsing systems (Gies 2003): $\sim 22 M_\odot$ for O8V stars and $\sim 18 M_\odot$ for O9V stars. Such values are very close to the minimum values found for HD 93161 A1 and A2, indicating a large inclination, most probably $>75^\circ$. With such an inclination, the system should undergo eclipses. Unfortunately, no photometric study of HD 93161 has been conducted up to now; additional observations are thus needed. They will help to constrain further the physical parameters of HD 93161 A.

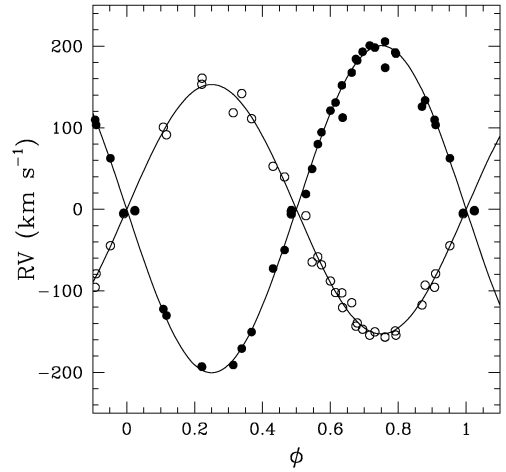


Figure 3. RV curves corresponding to the averaged RVs of HD 93161 A and assuming $e = 0.0$ (see Table 4). The open and filled circles correspond to the observed mean RVs of the primary and secondary, respectively. Note that we have used a reference frame with zero systemic velocity (see text).

6 XMM-NEWTON OBSERVATIONS OF HD 93161

6.1 Spectral analysis

Due to the proximity of HD 93160, the spectrum of HD 93161 was extracted from a circular region with a radius limited to 10 arcsec, as is shown in Fig. 4. Due to a poor calibration of the point spread function of the EPIC instruments (especially PN) at large off-axis angles, we extracted MOS and PN spectra in the extended source mode, i.e. without any EEF (encircled energy fraction) correction. For the chosen region, we estimate that this correction amounts to a factor ~ 1.8 . The background spectrum was obtained by selecting events from a circular region spatially offset from the source (see Fig. 4).

The X-ray data were then analysed with the XSPEC software (version 11.2.0). For each observation, the EPIC spectra were fitted simultaneously by a two-temperature thermal plasma MEKAL model

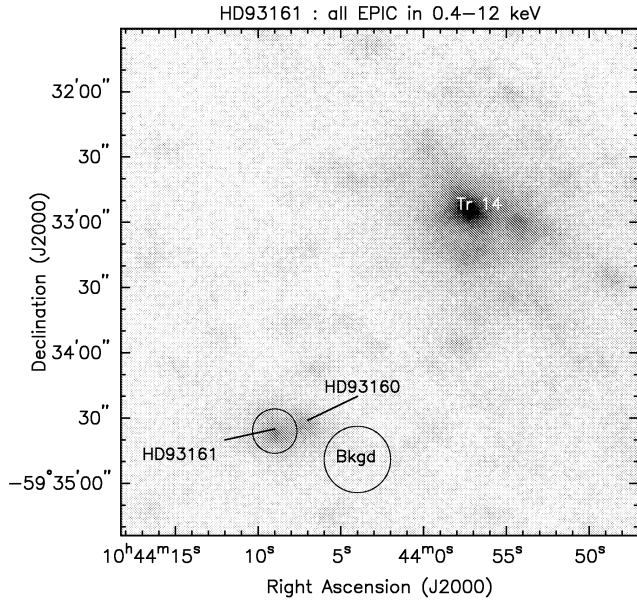


Figure 4. Combined EPIC image showing HD 93161, HD 93160 and their neighbourhood. The bright X-ray source to the top right is the core of the Tr 14 cluster, which contains notably Tr 14–1 and HD 93129AB. The source and background regions used to extract the X-ray spectrum of HD 93161 are also shown. Note that this image was exposure-corrected.

(Kaastra 1992; Mewe, Gronenschild & van den Oord 1985) of the form $\text{WABS}_{\text{int}}^* (\text{MEKAL}_1 + \text{WABS}^* \text{MEKAL}_2)$ where WABS_{int} was fixed to the interstellar absorption determined by Diplás & Savage (1994, $N_{\text{int}}^{\text{H}} = 4.5 \times 10^{21} \text{ cm}^{-2}$). The second WABS model is included to account for the possibility of additional absorption for the hard component of the spectrum. The abundances were set to solar. The best-fitting model parameters are presented in Table 5 and the spectra from revolution 284 are shown in Fig. 5 along with their best-fitting model. The spectral parameters and the fluxes are fairly stable throughout the five data sets. The parameters are quite typical for a single O-type star, with $kT_1 \sim 0.3 \text{ keV}$ along with a slightly higher temperature component. Our results are in agreement with the preliminary conclusions of Albacete Colombo et al. (2003) concerning the first two data sets.

6.2 Variability in the X-ray domain

We also investigated the variations of the X-ray emission of HD 93161, first on a short time-scale (within each pointing) and then on the long-term range (between the different pointings). For each

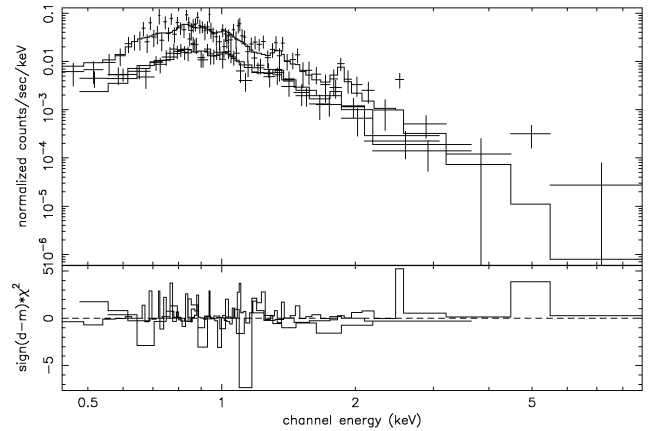


Figure 5. EPIC spectra of HD 93161 obtained for revolution 284. In the top panel, the upper and lower data correspond to the PN and MOS spectra, respectively, and the solid lines yield the best-fitting absorbed two-temperature model (see Table 5). The bottom panel shows the contributions of individual energy bins to the χ^2 of the fit. The contributions are carried over with the sign of the deviation (in the sense data minus model).

data set, the source and background events were extracted using the same regions as those defined for the spectrum analysis. After taking into account the flaring intervals (see Section 2.2), these event lists were used to construct background-subtracted light curves that were analysed by various methods (Sana et al. 2004). No significant short-term variability was consistently detected in the three instruments.

As a next step, we searched for variations between the different pointings. To this aim, we compared the mean count rates as well as the fluxes determined by spectral fitting (see Table 5); no significant long-term changes from one pointing to another were detected.

6.3 Wind interactions in HD 93161?

We have found no clear evidence that the components of HD 93161 are interacting in any manner. The visible spectrum presents only few emission lines and, in particular, no $\text{H}\alpha$ emission is seen that might arise in a wind collision zone (Thaller 1997; Sana, Rauw & Gosset 2001). Moreover, the minimum Roche lobe radii, $R_{\text{RL}} \sin i$, determined for the primary and secondary of HD 93161 A, are 24.1 ± 0.2 and 21.3 ± 0.2 , respectively (see Table 4). These large values contrast with the small typical radii ($< 10 R_{\odot}$) of O8–9V stars (Howarth & Prinja 1989; Papers II and III; Gies 2003; Sana et al. 2005b), indicating that this object is a well-detached binary system without any ongoing Roche lobe overflow. In addition, HD 93161 is not a particularly bright X-ray source, because its unabsorbed

Table 5. Best-fitting models and X-ray fluxes at Earth for each *XMM-Newton* observation of HD 93161. The fitted model has the form $\text{WABS}(N_{\text{int}}^{\text{H}})^* [\text{MEKAL}(kT_1) + \text{WABS}(N_{\text{int}}^{\text{H}})^* \text{MEKAL}(kT_2)]$, with $N_{\text{int}}^{\text{H}} = 4.5 \times 10^{21} \text{ cm}^{-2}$ (Diplás & Savage 1994). Quoted fluxes are in the 0.4–10. keV energy range and are not corrected for EEF (see Section 2.2). The unabsorbed fluxes $f_{\text{X}}^{\text{unabs}}$ are dereddened only considering the interstellar absorbing column. The quoted errors correspond to 1σ deviations.

Rev.	ϕ_{avg}	kT_1 (keV)	N_2^{H} (10^{22} cm^{-2})	kT_2 (keV)	$\chi^2_{\nu}(\text{dof})$	$f_{\text{X}}^{\text{abs}}$ ($10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$)	$f_{\text{X}}^{\text{unabs}}$ ($10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$)
115	0.18	0.26 ± 0.02	0.43 ± 0.13	0.71 ± 0.11	1.14 (132)	0.85	4.91
116	0.37	0.29 ± 0.03	0.44 ± 0.66	0.93 ± 0.67	1.26 (37)	0.77	4.46
283	0.17	0.27 ± 0.03	0.45 ± 0.14	0.67 ± 0.08	0.87 (106)	0.92	4.67
284	0.53	0.29 ± 0.02	0.44 ± 0.13	0.71 ± 0.12	0.89 (122)	0.96	4.72
285	0.75	0.29 ± 0.01	0.86 ± 0.16	0.79 ± 0.11	1.18 (139)	0.89	4.57

luminosity amounts only to $\sim 6.3 \times 10^{32}$ erg s $^{-1}$ (after EEF correction and for a distance of 2.5 kpc; see next section), or $\log(L_X/L_{\text{bol}}) = -6.6$ (using L_{bol} determined in the next section).

We can compare this observed X-ray luminosity to the emission expected for a wind collision in HD 93161 A. First, we determine the mass-loss rates \dot{M} and terminal velocities v_∞ using the recipes of Vink, de Koter & Lamers (2000, 2001) and the probable radii, masses, temperatures and bolometric luminosities of the stars (see Sections 5.2 and 7). Using the formalism of Pittard & Stevens (2002), we then find that, in the case of HD 93161 A, a radiative wind collision would lead to an X-ray emission with $L_X \sim 2 \times 10^{34}$ erg s $^{-1}$. However, the physical parameters of HD 93161 A (strengths of the winds, separation between the stars) should rather lead to a slightly adiabatic shock; the X-ray emission associated with such a collision could thus be significantly lower than in the radiative case. In fact, the total luminosity derived after applying the ‘canonical’ L_X-L_{bol} relation (see Berghöfer et al. 1997; Sana et al. 2005a) to the three massive stars of HD 93161 is $\sim 3.4 \times 10^{32}$ erg s $^{-1}$, indicating that the observed X-ray emission does not display a large excess attributable to a wind interaction. This conclusion is further supported by the lack of hard X-ray photons and the constancy of the X-ray emission. We might note however that our X-ray observations do not sample the whole orbital cycle; as was the case for CPD-41°7742 (Sana et al. 2005b), the signature of a wind collision could be limited to a few crucial phases.

Finally, Benaglia & Koribalski (2004) observed the Carina nebula with ATCA at 3 and 6 cm aiming at the detection of non-thermal radio emission, but HD 93161 was not detected in this survey, contrary to what would be expected if there was a wind collision between HD 93161 A and HD 93161 B. Therefore, from the available data, we conclude that there is probably no strong wind interaction in HD 93161.

7 FUNDAMENTAL PARAMETERS

Using the temperature scale of Martins, Schaerer & Hillier (2002), we can convert the spectral types derived above into effective temperatures of $38\,200 \pm 700$, $35\,250 \pm 700$ and $33\,700 \pm 800$ K for HD 93161 B, HD 93161 A1 and HD 93161 A2, respectively. The quoted uncertainties correspond to half a subclass for the components of HD 93161 A and to the range of temperatures for the O6.5V class according to Martins et al.

In order to derive the luminosities, we need the observed magnitude, the reddening, the brightness ratios between the different components, the bolometric corrections and an estimate of the distance. The WEBDA⁶ (Mermilliod 1988, 1992) data base lists the results of several photometric studies that provide the apparent V magnitude of HD 93161 (Feinstein, Marraco & Muzzio 1973; Walborn 1973; Forte 1978; Kaltcheva & Georgiev 1993; Massey & Johnson 1993). On average, we find $V = 7.84 \pm 0.02$ for the entire HD 93161 system. On the other hand, Vázquez et al. (1996) quote $V = 8.56$ and $V = 8.60$ for HD 93161 A and HD 93161 B, respectively. These magnitudes yield $V = 7.83$ for the entire system (components A + B) in good agreement with the mean value given above. Mason et al. (1998) found $V = 9.0$ for HD 93161 A and $V = 9.1$ for HD 93161 B. If these values are correct, the total magnitude of HD 93161 should be 8.3 in contradiction with the other observations. In the following, we therefore adopt apparent magnitudes of 8.56 ± 0.02 and 8.60 ± 0.02 for HD 93161 A and B, respectively.

⁶ Available at <http://obswww.unige.ch/webda/>.

The $(B-V)$ colour indices of HD 93161 A and HD 93161 B were found to be 0.20 and 0.23, respectively (Vázquez et al. 1996). Comparing with typical colour indices of O stars (Schmidt-Kaler 1982), we hence adopt $E(B-V) = 0.52$ and $E(B-V) = 0.55$ for HD 93161 A and HD 93161 B, respectively.

To estimate the relative brightness of the components of HD 93161 A in the visible spectral range, we evaluate the ratio of the EWs of the primary and secondary components for a number of lines in our 2004 FEROS spectra (the least contaminated by the B component). We use the He I $\lambda\lambda 4471, 5876$, He II $\lambda 4542$ and O III $\lambda 5592$ lines and we compare the EW ratios to the values for typical EWs evaluated from the compilation provided by Conti & Alschuler (1971) and Conti (1974). In this way, we estimate a brightness ratio of 1.7 ± 0.3 between the primary and secondary of HD 93161 A. This results in apparent magnitudes of 9.06 ± 0.07 and 9.64 ± 0.12 for the primary and secondary components, respectively.

The distances to the clusters Tr 14, Tr 16, Collinder 228 and Collinder 232 are still poorly known. How controversial the issue of the distances and even the actual existence of these clusters is, can be illustrated by a brief overview of the literature. Walborn (1995) summarized the situation as follows: Tr 14 is younger than Tr 16, but both clusters lie at the same distance; neither Collinder 228 nor 232 are genuine clusters, the former is part of Tr 16, whereas the latter contains stars of Tr 14 and Tr 16. The ambiguities on the distances of Tr 14 and Tr 16 are at least partly due to the uncertainties on the extinction towards these clusters. Using UBV photometry, Massey & Johnson (1993) inferred a distance of 3.2 kpc for Tr 14 and Tr 16. Assuming a normal extinction law and based on Strömgren and $H\beta$ photometry, Kaltcheva & Georgiev (1993) derived distances of individual stars in Tr 16. These authors reached a radically different conclusion: because our line of sight towards Tr 16 is almost parallel to the molecular cloud ridge in the Carina spiral arm, the stars might not belong to an actual cluster but could instead lie along an extended star formation region which is projected on a small area on the sky. Recently, Carraro et al. (2004) presented a study of Tr 14, Tr 16 and Collinder 232 based on $UBVR_I$ photometry. These authors argued that a unique reddening law is not appropriate to study the whole region and they estimated different values of the selective extinction $R_V = A_V/E(B-V)$ for each cluster. By comparing their observations with empirical main sequences, they derived distances of about 2.5 kpc for Tr 14 and Collinder 232, but a much larger distance of about 4 kpc for Tr 16. The latter result is clearly at odds with previous distance determinations. In particular, the studies of two eclipsing binaries in Tr 16 – Tr 16-1 (Freyhammer et al. 2001) and Tr 16-104 (Paper III) – have allowed us to obtain absolute luminosities of these stars. Adopting bolometric corrections from Humphreys & Mac Elroy (1984) and assuming $R_V = 3.3$, a distance of about 2.5 kpc is inferred for these binaries, in excellent agreement with the Tr 16 distance found by Tapia et al. (2003). In the following, we therefore adopt a distance of 2.5 kpc for HD 93161.

For Collinder 232, Carraro et al. (2004) infer $R_V = 3.73 \pm 0.03$, whereas they estimate $R_V = 3.83 \pm 0.33$ for HD 93161 and $R_V = 3.60 \pm 0.27$ for its close neighbour on the sky, HD 93160. Adopting an average value of $R_V = 3.72 \pm 0.16$, we obtain absolute magnitudes of -4.86 ± 0.11 , -4.29 ± 0.15 and -5.44 ± 0.09 for HD 93161 A1, A2 and B, respectively. Finally, adopting the bolometric corrections from Humphreys & Mac Elroy (1984), we obtain bolometric luminosities of $\log L_{\text{bol}}^{\text{A1}}/L_\odot = 5.25 \pm 0.04$, $\log L_{\text{bol}}^{\text{A2}}/L_\odot = 4.93 \pm 0.06$ and $\log L_{\text{bol}}^{\text{B}}/L_\odot = 5.53 \pm 0.04$.

Using these values, we have plotted in Fig. 6, the components of HD 93161 in the Hertzsprung–Russell diagram along with the

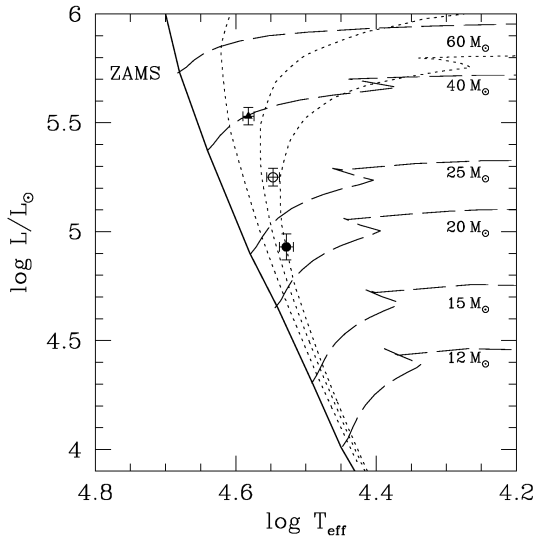


Figure 6. Location of components A1 (open circle), A2 (filled circle) and B (filled triangle) of HD 93161 in the Hertzsprung–Russell diagram. The evolutionary tracks are from Schaller et al. (1992). Isochrones corresponding to ages of 2, 3 and 4 Myr (dotted lines) are also sketched.

evolutionary tracks from Schaller et al. (1992) for solar metallicity. This diagram suggests that the stars have an age of 3–4 Myr and are thus slightly evolved off the zero-age main sequence (ZAMS), unlike other stars of the Tr 16 cluster, e.g. the components of Tr 16–104, which were found to lie very close to the ZAMS (Paper III). A crude interpolation between the evolutionary tracks yields evolutionary masses of 30.1 ± 1.1 and $22.4 \pm 0.8 M_{\odot}$ for the primary and secondary, respectively. This corresponds to a mass ratio of 1.34 ± 0.07 , in excellent agreement with the observed ratio of 1.31 ± 0.02 . The orbital inclination corresponding to these evolutionary masses would be $5^{\circ} \pm 2^{\circ}$, and this lower value of the inclination would prohibit eclipses. A photometric study of HD 93161 is thus necessary in order to assess the masses but we might already note that, in the case of the binary HD 93205 (Paper II), evolutionary masses are larger than the true masses. Finally, we caution that the luminosity of HD 93161 B might be overestimated if the star turned out to be another binary system.

8 CONCLUSIONS

We have presented a thorough spectroscopic investigation of HD 93161, based on medium- and high-resolution data obtained over a period of 7 yr. Although the small separation between the two components of HD 93161 did not enable us to obtain totally uncontaminated spectra, we have shown that HD 93161 actually harbours a detached SB2 system (HD 93161 A) and a star displaying RV variations (HD 93161 B).

The primary (secondary) star of HD 93161 A is most probably of spectral type O8V (O9V). Using the RVs measured for each component of HD 93161 A, we provide the first orbital solution for this system. The orbital period is 8.566 ± 0.004 d and the minimum masses for the primary and secondary stars are 22.2 ± 0.6 and $17.0 \pm 0.4 M_{\odot}$, respectively. Such large minimum values, compared to typical masses for O8V and O9V stars, suggest a high inclination of the orbit. It is therefore possible that HD 93161 A undergoes eclipses. A follow-up study should be undertaken to check the photometric behaviour of HD 93161 A.

The visual companion of HD 93161 A, HD 93161 B, is of type O6.5V(f) and is thus slightly hotter than its neighbours. HD 93161 B presents RV variations, but no fully obvious signature of a companion could be found in its spectrum. The small number of our observations of this star did not enable us to investigate the periodicity of the velocity changes. Additional observations are required to check the binary status of HD 93161 B.

Considering that HD 93161 lies at a distance of 2.5 kpc, we found that all three components of HD 93161 are slightly evolved off the ZAMS. In the X-ray domain, HD 93161 appears as a moderately bright X-ray source displaying a two-temperature spectrum and no significant variability, suggesting an X-ray emission predominantly intrinsic to the stellar winds of the individual components.

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